Jets

- experimental aspects

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Contents

- What is a jet and how to measure it? Relation to theory and partons (1)
- Heavy-ion collisions experimental controls some theoretical expectations... (1->2)
- Probing quark-gluon plasma with jets and jetlike observables; jet-medium interactions (2+3)
- All this based on: lessons from RHIC & LHC
	- Note: choice of experiments is randomized much of what shown measured by multiple experiments...

Thanks to all the authors/experiments for the graphics/slides shamelessly stolen for the purpose of this talk

What is a jet?

A spray of collimated showers/particles - Hardly ever better defined...

Jet = Parton AND its radiation

Note: experiment measures spray of particles (~hadrons)

Jets (unlike single hadrons) are objects which are "better" understood/calculable within pQCD

S.D Drell, D.J.Levy and T.M. Yan, Phys. Rev. **187**, 2159 (1969) N. Cabibbo, G. Parisi and M. Testa, Lett. Nuovo Cimento **4**,35 (1970) J.D. Bjorken and S.D. Brodsky, Phys. Rev. D 1, 1416 (1970) Sterman and Weinberg, Phys. Rev. Lett. 39, 1436 (1977) ...

Jet finding

Note: jets originate from hard partons, however definition of a parton in terms of a jet is ambiguous -> multiple jet definitions.

Optimum jet finder algorithm

Several important properties that should be met by a jet definition are $[3]$:

1. Simple to implement in an experimental analysis;

2. Simple to implement in the theoretical calculation;

3. Defined at any order of perturbation theory;

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4. Yields finite cross section at any order of perturbation theory;

5. Yields a cross section that is relatively insensitive to hadronization.

Tevatron 1990

... and infrared safe and colinear safe (~2000)

GCD divergencies and jet finders

QCD probability for gluon bremsstrahlung at angle θ and \bot -mom. k_t :

$$
dP \propto \alpha_s \frac{d\theta}{\theta} \frac{dk_t}{k_t}
$$

For pQCD to make sense, the (hard) jets should not change when

• one has a collinear splitting *i.e.* replaces one parton by two at the same place (η, ϕ)

• one has a soft emission *i.e.* adds a very soft gluon

Jet algorithms:

Colinear & infra-red safety

Safety: Results = jets = reconstructed objects - insensitive to modifications at the soft scale of radiation (hadronization, soft colin. radiation)

Collinear safety

Infrared safety

Modern jet algorithms

- Colinear and infrared safe
- Improved performance

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- Rigorous definition of jet area
- Different algorithms -> different response to the underlying event
	- •Developed for uniform bg subtraction (pile-up) at LHC

Two main classes of algorithms: recombination (kt, Cambridge/Aachen, anti-kt) and cone (Mid point cone, CDF, SIScone)

Sequential recombination (clustering) algorithms

Majority of QCD branching is soft & collinear, with following divergences:

$$
[dk_j||M_{g\to g_ig_j}^2(k_j)] \simeq \frac{2\alpha_s C_A}{\pi} \frac{dE_j}{\min(E_i,E_j)} \frac{d\theta_{ij}}{\theta_{ij}}, \qquad (E_j \ll E_i, \ \theta_{ij} \ll 1).
$$

To invert branching process, take pair with strongest divergence between them - they're the most likely to belong together.

This is basis of k_t/D urham algorithm (e^+e^-) :

1. Calculate (or update) distances between all particles *i* and *j*:

$$
y_{ij}=\frac{2\min(E_i^2,E_j^2)(1-\cos\theta_{ij})}{Q^2}
$$

2. Find smallest of y_{ii}

NB: relative k_t between particles

- If $>$ y_{cut} , stop clustering
- \triangleright Otherwise recombine *i* and *j*, and repeat from step 1

Catani, Dokshitzer, Olsson, Turnock & Webber '91

Example: Kr algorithm

k_t jet algorithm 1.1

The definition of the inclusive k_t jet algorithm that is coded is as follows:

1. For each pair of particles i, j work out the k_t distance

$$
d_{ij} = \min(k_{ti}^2, k_{tj}^2) \,\Delta R_{ij}^2 / R^2 \tag{1}
$$

with $\Delta R_{ii}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$, where k_{ti} , y_i and ϕ_i are the transverse momentum, rapidity and azimuth of particle i and R is a jet-radius parameter usually taken of order 1; for each parton *i* also work out the beam distance $d_{iB} = k_{ii}^2$.

- 2. Find the minimum d_{\min} of all the d_{ij}, d_{iB} . If d_{\min} is a d_{ij} merge particles i and j into a single particle, summing their four-momenta (this is E-scheme recombination); if it is a d_{iB} then declare particle i to be a final jet and remove it from the list.
- 3. Repeat from step 1 until no particles are left.

Anti-kt: k_t^2 is replaced by k_t^{-1}

M. Cacciari, G. P. Salam, G. Soyez JHEP 0804:063,2008. e-Print: arXiv:0802.1189 [hep-ph]

Cone algorithms

Jet Cones

- Cones are always understood as circles in rapidity (y) and azimuth φ.
- A particle i is within the cone of radius R around the axis a if

– Find directions of dominant energy flow " find \sim

- $\Delta R^2_{ia} = (y_i y_a)^2 + (\phi_i \phi_a)^2 < R^2$
- ... usual hadron collider variables
- $-$ Typical: R = 0.4 0.7

Basic Idea: Two contracts of algorithms of algorithm

– center of the cone ≡ direction of the total momentum of its particle contents $\overline{}$ $\overline{\$

Speed matters!

FJ: Significant gain for high-multiplicities

Jet finding - jet finders

Complete suite of algorithms - FastJet package:<http://www.lpthe.jussieu.fr/~salam/fastjet/>

Jet shape - R-dependence

 $\frac{1}{2}$

Jets in collider experiments

Hadronic collisions: pQCD and jets

Inclusive jet production: pQCD & data

Jets are fairly well known by now... and well described by theory and MC => attractive tool for heary-ions

JET composition Particle Flow Jet Reconstruction Particle Flow ⁵ TotaLlet Energy Correction Factors are fraction of hadronic energy observed in reconstructed jets is currently overestimated, both in ruction[.] 21 14

 $C(\overline{p}_T^{raw}, \eta) = C_{MCtruth} (p_T^{raw}, \eta) \times C_{Residual} (p_T^{raw} \cdot C_{MCtruth} (p_T^{raw}, \eta)$,

The overall jet energy correction of the and its ancertainty unshown in Fig. 1 of for fixed jet *p*_I values. C_c ALO jets require a much arger correction factor factors of the correction of cor pared to the track-based algorithms. In the region beyond the tracker cover are in agreement within the state of the systematic uncertainties. **Figure 13 shows the correction factors and i** the r uncertainty as a function of the system of the ixed η values. The systematic uncertainty as a function of the systematic uncertainty as a function of the systematic uncertainty as a function of the systematic unc overall calibration factor is the sum in the sum in the relative scale and the absolute scale and the abso un extainties. Figure 14 shows the combined uncertainty of the jet energy sca fur then of jet p_T while Fig. 15 shows the same quantity as a function of η . **165% charged hadrons** the three and uncentairy efectrions for fixed jet an yalues (GRNGE) SIGHB)

Jy Correction Factor 24e L*unknown* 1 rest trom 2H 2 Veed to 3 Calorimeter jets JET-POLLSET rack gever to be to be 7 Jerus Particle Flow jets anti-k_T R = 0.5 \rightarrow \sim 2.5 Measure SDGeV Know" Measure a jet? y Correction Factor 3 Calorimeter jets Need to Verpoverrack published to the all component Rus-Track jets **Particle Flow jets** \sqrt{s} anti-l $p_T =$ **CMSF Total JES Correction Factor** $the LyRR$ pown] rest from $BATEA + MC$

Jet: from particular to the particular of **the physics event generation** 9:;<3=< 5,>+=?<

 $J(\overrightarrow{p}_{partons}) \approx J(\overrightarrow{p}_{shower}) \approx J(\overrightarrow{p}_{hadrons}) \approx J(\overrightarrow{p}_{cells/tracks})$

Jets

A Jet Detector

 $\overline{\text{MS}}$

Improvements in jet reconstruction on detector level => Particle flow

Jet: energy scale &

resolution

Control over the two crucial in p-p and AA collisions

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Width == Resolution

JET: From Measured to

 A_{α} and A_{α} This is an experimental enterprise! It is a substantial effort...

ATLAS Linearity with, data Control of the energy scale - ATLAS - linearity

Jet Resolution Measurement Solution

An example: proton-proton collisions

Heary-ion collisions

RYIC to LYC

The hot-gCD laboratories

The Relativistic Heavy Ion Collider (BNL)

The hot-QCD laboratories **The-hot-QCD-laboratories-**

QCD phase diagram

'heoretical landsca - theoretical landscape

As hot as an early universe? Primodal soup of quarks in the laboratory?

Some history...

Hot QCD Labs

Strategy: how to study QCD matter experimentally?

- **Need to find those observables that:**
	- $-$ Are sensitive to crucial parameters of hot QCD matter
	- $-$ Can be modeled well theoretical understanding
	- **Can\$be\$measured\$well\$–\$experimental\$control\$**
	- $-$ Can connect theory and data
- => Inclusive measurements; correlations; compare with more elementary collisions (p-p, p-A); compare different energy regimes

Heavy-ion collisions

Collision evolution

Note: hard scatterings occur early (at t~0)! High energy partons "witness" the evolution and jets "testify" about their fate/CV

evolution

We are interested in properties of the QGP phase Need to disentangle effects from different phases - not a simple problem by principle: detectors do NOT measure these time-periods/phases separately => need for detail understanding of the physics processes, particle production, dynamics of the system in each phase(!), etc => modeling, various assumptions may play an important role in physics interpretation Need for control of the initial conditions, geometry of the collision, the incoming parton distributions (nuclear-PDF vs nucleon-PDF) ...

HI collisions: Particle production

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HI collisions: Particle production

non-singly diffractions as a function of √sNN. Curves and √sNN. Curves are fitted to the two data sets. Right o

Feedback within the heavy-ion community:

- **1. Multiplicity is crucial [input] for modeling** θ by similar theoretical approaches separated by dashed lines. See text and Ref. θ \parallel **T.** IVIUILI
	- **2. Saturation models tend to predict lower multiplicity**

3. Data driven extrapolations did not seem to anticipate the **results#** α = 168 (a) α (sign.) α β (see also α). This is shown is shown in Fig. 2 (right panel) with panel) α predictions from various models. As a whole the perturbative \mathbf{r} whole the perturbative \mathbf{r}

Carlo models (figure, notation and references used in Fig. 2 are from Ref. [3]) based

Systematic control: RHIC vs LHC

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End of 1/3

Until now...

- Jets in elementary collisions: must specify an operational definition (algorithm, R, recombination scheme); variety of infrared and collinear safe algorithms
- Jet measurements in e-e and p-p collisions under control - experimental and theoretical understanding - although proper jet reconstruction is an effort even in the "simple" case (vacuum)
- HI collisions: hot QCD matter; large particle (production) densities as compared to vacuum - evolving with centrality

... back to "calibration" measurements

HI collisions: Particle production

predictions from various models. As a whole the perturbative QCD-inspired Monte

Particle production: source

dimensions ⁹

- The systematics of the product of the three radii is shown in Fig. 4. The product of the product of the radii is shown in Fig. 4. The product of the radii is shown in Fig. 4. The product of the radii, which is shown in Fig **•** system with larger (2x) volume and (1.4x) lifetime (w.r.t RHIC); follows the **trend of multiplicity; faster expansion <=> larger collective flow 2. Pair momentum dependence:** \blacksquare **1. Energy dependence:** Phys.Lett.B 696:328-337,2011
- $\overline{}$ is inverse way. The size of the size of the size of the velocity gradient of $\overline{}$ **•** larger radii, strong dependence on kT; Rout/Rside smaller than at RHIC; **Example 22. The magnitude of** *Rlooverall agreement with extrapolations* **of the total duration of the total dura**
- \vert 2 lmnortant constrains to [hydrodynamical] modelling $f(x) = \frac{1}{\sqrt{2\pi}} \int_{0}^{x} \frac{dx}{\sqrt{2\pi}} dx$ **3. Important constrains to [hydrodynamical] modelling**

to those obtained for central gold and lead collisions at lower energies at the AGS [35], SPS [36, 37, 38], and

RHIC [39, 40, 41, 42, 30, 43].

Particle production: source

dimensions 9

- obtained for central gold and lead collisions at lower energies at the AGS [35], SPS [36, 37, 38], and RHIC [39, piume and **Example 10 and 10 start of multiplicity; faster expansion <=> larger collective flow** $t_{\rm c}$ and $t_{\rm s}$ and lead collisions at the AGS \sim **•** system with larger (2x) volume and (1.4x) lifetime (w.r.t RHIC); follows the
- for midrapidity pions exceeds 10 fm/*c* which is 40% larger than at RHIC. These results, taken together **2. Pair momentum dependence:** \blacksquare pseudorapidity density and is two times larger at the LHC than at RHIC.

time \$ *^f* can be obtained by fitting *R*long with

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longitudinal expansion, i.e. to the decoupling time of the system [31]. Quantitatively, the decoupling

 θ ones on LT , $\mathsf{Point}(23, 34)$ anology then at $\mathsf{DUL}(2, 11)$ that the first formulations at \mathbf{u}_i is hotter, and \mathbf{u}_i is a larger structure, and \mathbf{u}_i **example 20 as conserved to lower energies. The size of the versely proportions** • larger radii, strong dependence on kT; Rout/Rside smaller than at RHIC;

The ALICE Collaboration would like to thank all its engineers and technicians for the technicians for the their

Acknowledgements $\frac{1}{2}$ become who who constructive to flex due developed in a decrease $\frac{1}{2}$ **12. Important constrains to [hydrodynamical] modelling**

Identified particles T_{\perp} at C_{\perp} and I_{\perp} boenines parties Identified particles 0-5% central collisions. Results from STAR and PHENIX are also shown for Au–Au

& expansion of the system R average in R the C 15 to R expansion of the system

 \vert ALICE: excellent particle identification capabilities at the LHC \vert temperature and the average radial boost velocity from the average radial boost velocity from the Blast Wave m
[7] In the Blast Wave model.[7] In the Blast Wave model.[7] In the Blast Wave model.[7] In the Blast Wave mode

temperature and the average radial boost velocity from the Blast Wave model.[7]

Statistical hadronization of the

All yields (but protons) described by thermal model with **T_{ch}=164 MeV (and** μ_b **=1 MeV)**

- Similar temperature as at RHIC, however proton/pion below the fit the tension already procent at PHIC $\frac{1}{2}$ similar temperature as a set $\frac{1}{2}$ for $\frac{1}{2}$ already present at RHIC and the second by the second for
- Strange particles constrain fit
- $\frac{1}{2}$ sions are modernidependent (commun • Conclusions are model independent (confirmed with THERMUS) Strange particles constrain fit
Conclusions are model independent (confirmed with THFRMUS) ● Similar temperature as at RHIC, however proton/pion below fit

● Strange particles constrain fit

 $-$ Tension already present at RHIC and α

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Experimental control of collision

Centrality measurement: use of Glauber in an experiment

- Fraction of cross section, 2 approaches:
	- Fit with Glauber Monte Carlo
	- Correct: subtract BG, efficiency and integrate multiplicity distributions
- N_{part} , N_{coll} , N_{spect} : require Glauber fit (computed using cuts on impact parameter)
- Estimators:

V0, SPD clusters, TPC tracks, ZDCs, …

• ZDC measures N_{spect} : test of Glauber picture

- Glauber fit ingredients
	- Woods-Saxon (constrained by low energy electron-nucleus scattering)
	- Inelastic pp cross section (measured by ALICE)
	- Nucleons follow straight line trajectories, interact based on their distance
	- Compute (fit) observables assuming: $N_{\text{ancestors}} = \alpha \cdot N_{\text{part}} + (1 - \alpha) \cdot N_{\text{coll}}$

only charged particles visible

Peripheral Collision

 $Color \Rightarrow Energy loss$ in TPC gas

only charged particles visible

Central Collision

200 GeV Au+Au: **Nch~4800**

Nuclear geometry - Glauber model

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phase space gives a total transverse energy per pair of participating nucleons of 92 *±* 6 GeV for Energy density: RHIC to LHC

on h*N*parti by normalizing d*E*T/d*h* by the number of participating pairs of nucleons, h*N*parti/2. $LHC > 2.5 \times R4IC$

... within a volume (per nucleon)

Figure 1: Transverse energy density versus *|h|* distribution for a range of centralities of (0– $Vert hot.5\mu per dense(?)$ -> how The statistical uncertainties are negligible. Also shown are a Gaussian fit and the predictions \mathcal{A} Figure 2: Transverse energy density normalized by (h*N*parti/2) versus h*N*parti for PbPb colsuper dense(?) -> how dense...? certainties. The statistical uncertainties are negligible. Lower energy PHENIX data are also Very hot, super dense(?) -> how dense...?

Hadron production in $\sqrt{10}$ dhom ahodiightionin Hadron production in

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heavy-ion collisions ^s ratio increases and develops a maximum, reaching a ratio Λ/ K⁰ ~ 1.5 for pT ~ 3-3.5 GeV/c in 0-5% central collisions. A comparison with resultsa from RHIC for 0-5% central and 60-80% peripheral Au-Au collisions in Fig. 5 (right panel)

A slight digression... s singlet \mathcal{R} HIC vs LHC

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(LHC: higher mean pT - more flow)

Much more baryons than mesons in central for the minimum bills of minimum bills collisions as compared to proton-proton and periodic ratios from the left collisions as compared to proton-proton panel see text paralise in similar at slight rations at Slight (coalescence/recombination? bulk+jet?)

LHC similar to RHIC Maximum at slightly higher-pT

 t hermal e_{α} and e_{α} and e_{α} and peripheral collision ratios from the left peripheral collis thermal

 e_{α} and e_{α} and e_{α} and peripheral collision ratios from the left peripheral collis

Hadronization of bulk+hard

- parton coalescence - parcon coalescence

Probing the

unknown medium...

QED: Passage of electrically

charged particle through matter h

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Brehmsstrahlung in QCD:

Formation time -> coherence effects time -> coherence

Brehmsstrahlung in QCD **Bremsstrahlung-energy-loss-in-QCD**λ<τ=> Multiple scatterings add coherently $t_{formation} < L \Leftrightarrow \omega < \omega_c$ High(energy(**color-charged-probe** propagating through color charged medium ω =xE (LPM effect; multiple soft radiations) 100000000000000

Partonic energy loss in QCD medium is proportional:

Production

Hard

- to squared average path length (Note: QED ~ linear)
- to density of the medium

Define a transport coefficient:

 $\hat{q} \sim \mu^2/\lambda$

⇒ energy flow (parton+radiation) modified as compared to jet in vacuum

 $-dE/dx \sim \alpha_s \hat q L^2$

 ρ

 ω =(1-x)E

 $\lambda \propto \frac{1}{\epsilon}$

⇒ jet "quenched" ("softened" fragmentation) 22)
22(QCD(laboratory(with(heavy9ions,(MPloskonn)
22(QCD)

Generic expectations from energy loss

Longitudinal modification: out-of-cone: energy lost, loss of yield, di-jet energy imbalance in-cone: softening of fragmentation

Transverse modification

out-of-cone: increase acoplanarity kt in-cone: broadening of jet-profile

Factorization in heavy-ion collisions? 67

Jets in heavy-ion collisions

- an idealization **Jets-in-heavy-ion-collisions:-**

Jets in heavy-ion collisions RHIC & LHC

LHC + RHIC: QCD evolution of jet quenching ?

Vary energy of the jet: LHC: Vary the scale with which QGP is probed (a la DIS) Compare and contrast RHIC and LHC

Jets in HI collisions & Experimental difficulties: Vacuum jet vs jet on top of the HI background...

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Jets in heavy-ion collisions RHIC & LHC

Jets in heavy-ion environment - few experimental notes:

- large combinatorial backgrounds (especially at RHIC) - energy within an event varies from point to point ("fluctuations") - a plus for LHC is larger kinematic reach - abundance of highenergy jets (higher-pT measurements less affected by backgrounds) => various approaches among experiments for background suppression AND/OR jet energy-resolution corrections - is there an optimal jet definition for heavy-ion collisions (?) => use multiple jet algorithms (?); sub-jets (?); filtering (?) - jets are reported on the particle (generator) level - hadronization corrections (to the "parton" jet) in HI collisions impossible

"Easier" (than full jet reconstruction) exercise: Jet-quenching via leading hadrons

Measured as a function of collision centrality

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Hadron suppression

 $R_{AB} = \frac{d^2N/dp_t d\eta}{T_{AB}d^2\sigma^{pp}/dp_t d_\eta}$ $T_{AB} = \langle N_{bin} > \langle \sigma_{inel}^{pp} \rangle$

N N

Nuclear modification factor:

#(particles observed in AA collision per N-N (binary) collision)

7 #(particles observed per p-p collision)

 α' algoritical (205 is far $RAA = 1$ of bight "No effect" case is for $RAA = 1$ at high pT where hard processes dominate

 KAA

Jet quenching - RHIC

RAA: extreme scenarios

Note: I am not showing you the $P(\Delta E)$

*P(*Δ*E) - probability for parton to loose* Δ*E*

Brehmsstrahlung in QCD **Bremsstrahlung-energy-loss-in-QCD-**

 $t_{formation} < L \Leftrightarrow \omega < \omega_c$

λ<τ => Multiple scatterings add coherently

 $dx \sim \alpha_s qL$

 ρ

 $\lambda \propto \frac{1}{\epsilon}$

High energy color charged probe propagating through color charged medium

-> different pollisions systems? An idea: vary the path length experimentally? -> sensitivity to the collision profile

Partonic energy loss in $Q\mathcal{D}$ medium is proportional:

- to squared average path length (Note: QED ~ linear)
- to density of the medium

⇠ *^µ*²*/*

q ˆ

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Azimuthal angular asymmetry in particle production

 $\overline{}$ Azimuthal ar Azimuthal anisotropy

APS Viewpoint: A "Little Bang" arrives at the LHC (E. Shuryak)

- **1. Collective behavior observed in Pb-Pb collisions at LHC (integrated:** +0.3 v_2^{RHIC} – consequence of larger <p_T>) -> $v_2(p_T)$ similar to RHIC – almost ideal fluid at LHC ? Similar observation down to 39GeV!
- **2. New input to the energy dependence of collective flow**
- **3. Additional constraints on Eq-Of-State and transport properties**

 \mathbf{L} out-of-plane

Suppression out-of-plane stronger <= longer in-medium path length - significant effect even at 20 GeV/c **Path length dependence of energy loss** ? Additional constraints to energy loss models (?)

- similar information from v2 at high n - similar information from v2 at high p_T $R = \frac{1}{2}$

$$
R_{AA}(\varphi) = R_{AA}(1 + 2v_2 \cos 2(\varphi - \psi))
$$

RAA for different particle type

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 $\Delta E \propto \alpha_S C_R q L^2$

3 for gluons

Is parton energy loss different for gluons, light-quarks and heavy-quarks?

Expectation: $\Delta \epsilon_g > \Delta \epsilon_{light-g} > \Delta \epsilon_{heavy-g}$ Casimir (color factor) - gluons "glue" better to the medium than quarks "Dead-cone" effect: mass of the parent quark => radiation for angles θ <m/E $CR = 4/3$ for quarks,

is suppressed

=> RAApions < RAAD-mesons < RAAB-mesons

RAA for different particle type

Discussion based on LHC results

p_T (GeV/c) pendenge of 40 access on **AA** $2)$ Weak dependence of 40° QC $^{20}_{60}$ e-loss on 0.2 0.4 0.6 0.8 $LangBCSK0$ R and delow $7-GeV$ - $1.2 \leftarrow$ J/ψ fram B -decays $\frac{1}{\sqrt{2}}$ dead cone vertices 1.6 e-loss; **V&E**s dep. on mass?; Changed Darticle 0-5% (CMS) 2 **= 2.76 TeV NN PbPb s** $factor$? 1.6 5 mol eff cf fc t ?) \rightarrow Charged Particle 0 - 5% (ALICE) **ISOLATED Photon 0 - 10% (CMS) Theory Ch. hadron (YaJEM-D) Ch. hadron (HT-M) Ch. hadron (JEWEL) Ch. hadron (HT-W) Ch. hadron (GLV)** Similar suppression for heavier-q (strange, etharm) and gluons (large elastic manifestation of flow (?) Rise towards higher 1) Harder partonic spectrum (as compared to RHIC) parton energy

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RAA for different particle type

Discussion based on LHC results

Similar suppression for heavier-q (strange, charm) and gluons (large elastic e-loss; less dep. on mass?; color factor? - small effect?) J/ψ from B-decays - dead cone effect?

Lambda vs K0 RAA below 7 GeV – manifestation of flow (?)

Rise towards higher pT's: 1) Harder partonic spectrum (as compared to RHIC) 2) Weak dependence of [pQCD] e-loss on parton energy

Photons and Z's not suppressed -> quenching is a final state effect

Heavy-flavor suppression

 \parallel Variants of radiative++ energy loss agree with data

Until now...

- Jets in elementary collisions: must specify an operational definition (algorithm, R, recombination scheme); variety of infrared and collinear safe algorithms; under control theory/experiment;
- HI collisions: hot QCD matter; large particle (production) densities as compared to vacuum - evolving with centrality; Jet measurements difficult (Today you will see that possible nevertheless)
- Leading hadrons suppressed <-> parton energy loss (jet quenching); Hadrons select particular ensemble of jets(!) - fragmentation bias (more Today) - relation of parton vs hadron energy (?)

... back to jet quenching measurements

"Easier" (than full jet reconstruction) exercise: Jet-quenching via leading hadrons

Sensitivity of particle correlations to different underlying physics

Two-particle correlations

Conditional yields - LHC

Yield per trigger particle AA/pp-> IAA (unity==no effect) Yield per trigger particle AA/pp-> IAA (unity==no effect)

PRL 108, 092301 (2012)

 $\overline{2}$ – Near*side"enhanced"(~1.2)"!"Change"in"fragmenta:on"func:on,"Change"in"quark/gluon"jet"ra:o," •Central events:

= 2.76 TeV NN *s*

 $\frac{e}{b}$ 1.0 $\frac{1}{\sqrt{1-\frac{1}{n}}}\int_{0}^{\frac{1}{n}} P(x)dx$ $\frac{1}{2}$ bkg ² *v* • near-side enhancement (>1: change in FF? bias on parton spectrum?; g/g-mix $\frac{1}{2}$ \mathcal{L} different in PbPb as compared to p-p?) - consistent with jet quenching... • recoil: suppressed - consistent with quenching

IAA: data & theory description ⁹² ta & theory desc

$\overline{}$ Near-side enhancement: | Near-side enhancement:

- \vert \vert \vert Reproduced by AdS/CFT inspired (L³ path length dependence) and ASW inspired (L²) models \vert $-$ Reproduced by AdS/CFT - inspired (L³ path length dependence) and ASW - inspired (L²) models
- \vert $-$ YaJEM too high (L dependence)

 \mathcal{L} and \mathcal{L} – Reproduced"by"all"except"YaJEM"" μ and μ ok, so... compatible with jet guenching...

RHIC Example: High-pT hadrons - quantitative analysis

Model calculation: ASW quenching weights, detailed geometry Simultaneous fit to data.

Reasonably self-consistent fit of independent observables Main limitation is the accuracy of the theory...

So, why bother with full jet reconstruction in heavy-ion collisions?

RAA and correlations of leading hadrons provide constraints on density of the medium (ghat), however do not tell us about the \ast parton \ast energy loss and its dynamics; leading hadrons are biased towards jets that interact little or not at all with the medium

So called surface bias: requesting a high-pT particle selects a population of jets close to surface of the medium these jets interact only little (or not at all) with the medium

=> full jet reconstruction premise: integrate over the hadronic degrees of freedom; better access to the parton energy scale; dynamics of the jet quenching (?); other promising observables: gamma-jet correlations

Jets in heavy-ion collisions

LHC + RHIC: QCD evolution of jet quenching ?

Vary energy of the jet: LHC: Vary the scale with which QGP is probed (a la DIS) Compare and contrast RHIC and LHC

Jets in HI collisions & Experimental difficulties: Vacuum jet vs jet on top of the HI background...

Jets in HI collisions & Experimental difficulties: Vacuum jet vs jet on top of the HI background...

HI jet finding:

Many of these objects are simply background $W = \frac{1}{2} \int_{-\infty}^{\infty} \frac{1}{2} \left(\frac{1}{2} \right) \left(\frac{1}{2} \right) \left(\frac{1}{2} \right) \left(\frac{1}{2} \right)$ A single event: all particles clustered ("assigned") to a jet Energy of the signal jets overestimated due to background energy => several possibilities to subtract the average background and/or suppress the background particles Land background jets]

Background subtraction

Developed for pile-up rejection in p-p....
\n
$$
p_T^{jet} = p_T^{cluster} - \rho \times Area
$$
\n
$$
p_T^{jet} = p_T^{true} \otimes \delta \rho
$$

- $ρ:$ median pT per unit area of the diffuse background in an event $$ measured using background "jets" as found by kT algorithm
- A: area of the jet measured using number of artificially injected infinitely soft particles of finite "size" into an event that are clustered into the jet
- $\delta \rho$: uncertainty due to noise fluctuations $-$ non-uniformity of the event background

M. Cacciari, G.Salam Phys.Lett.B659:119-126,2008. e-Print: arXiv:0707.1378 [hep-ph]

Must correct for remaining residual energy resolution - magnitude of the correction is related to the background fluctuations - jet Area : small R (area) - smaller correction

In Jet reconstruction in HI collisions: Background fluctuations: characterized by δ pt; spectrum before corrections Schule: Only On My 4000, 200 Inhomogeneous structure of events. Background fluct measured Pb-Pb events.

 $\overline{2}$ step procedure to correct for $\overline{2}$ 101

Energy resolution deteriorated due to background energy fluctuations

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Background corrections 103

in Atlas **Jet Assets** (fakes)

- Reconstruction algorithm anti- k_{t} (0.2, 0.4). з,
- Input: calorimeter towers 0.1×0.1 (Δη \times Δφ). 鱼
- Event-by-event background subtraction: $E_{\text{Tsub}}^{\text{cell}} = E_{\text{T}}^{\text{cell}} \rho^{\text{layer}}(\eta) \times A^{\text{cell}}$
- Anti-k_, reconstruction prior to a background subtraction.
- ▶ Underlying event estimated for each longitudinal layer and n slice separately.
- We exclude jets with $D \! = \! E_{\mathit{T\,tower}}^{\mathit{max}} / \langle\mathit{E_{\mathit{T\,tower}}}\rangle \! > \! 4$ to avoid biasing subtraction from jets but no jet rejection based on *D.*
- Iteration step to exclude jets with E_{τ} > 50 GeV from background estimation. Jets corrected for flow contribution.

- **UE fluctuations from soft particles can be reconstructed as jets (fakes)**
	- Worse for larger *R*, contribute up to ~80 GeV
	- Require additional signal of **hard particle** production
	- Reject fakes by requiring jet to match:
		- Track jets or EM clusters with $p_T > 7$ GeV
	- Residual fake rate estimated to be ~3% at 50 GeV

Background subtraction / jet energy

corrections (CMS)

PF pseudo-tower

η strip

$\begin{bmatrix} x \\ y \end{bmatrix}$ are constructed by the construction of $\begin{bmatrix} x \\ y \end{bmatrix}$ $\mathcal{O}_\mathcal{A}$ grid according to HCAL cell dimensions $\mathcal{O}_\mathcal{A}$ a) Event-by-event subtraction of the heavy-ion background

experience and dispersions and dispersion and dispersions are the dispersion of the set of - Reconstructed particles towered into an (η,φ) grid according to HCAL cell dimensions

- calculated for example and dis-• Mean tower energy and dispersion are calculated for each η strip
- Same iterative background subtraction Same iterative background subtraction applied in [0], described in [1]
- andom cone studies. good agreed in the service in the service of the service of the service of the service of
And the service of t • Random cone studies: good agreement between background fluctuations in data and HYDJET simulations

 \cdot The effect of quenching on the energy scale is constrained using the associated charged particle spectra jet associated charged particle spectra

HYDJET simulations simulation of PYTHIA jets b) Jet energy corrections (JEC) based on GEANT

 \mathbf{v} is constrained using the jet associated using the jet associated using the jet associated using \mathbf{v} charged particle spectral in LIVP c) Validation of the BG subtraction + JEC for PYTHIA jets embedded in HYDJET

• Jet quenching measurements with fully reconstructed jets

105
Jet R central-Peripheral 60-80%) Ω Can \mathcal{I} · *Scentral-Peripheral 60-80%*)

 \mathcal{L} Medium-induced radiation can distribute \mathcal{L}

RCP : similar as RAA, but denominator are not yields from proton-proton but from peripheral heavy-ion collisions proton-proton but from peripheral heavy-ion collisions as KAA, but denominator ‣ Can gain additional insight by considering **inclusive** energy loss

80 Rcp ~ 0.5 => suppression - jets loose energy in most central events - the radiation is not captured within the jet cone (R) solute scale this additional suppression is not large, it is more significant in comparison to the *Z*₀ or Delle-Yan production provided provi cesses [17, 32, 35]. These latter channels are dominated by $q_{\text{max}}(P)$ arises primarily from *g*+*g* (and *g*+*q*(*q*) at larger *ET*) processes.

106

t,

Suppression as a function of R **Quantitative Statement of** *R* **Dependence**

Jet quenching with charged jets -

ALICE

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Another observable: Ratio of x-sections: $R_{1}< R_{2}$ P, IP do P, IP $f(x)$ for $f(x)$ and $f(x)$ is $f(x)$ R1/R2 where R1<R2 $RVR2$ where $R15R2$

 $NOTE: Systematically different$ **Ratically different** Ratio measurement - same effect found! 05/31/12 ¹⁴ **R=0.3**

 f' peripheral and collisions of f' Ratio R=0.2/R=0.3 consistent with vacuum jets for peripheral and central collisions

05/31/12 13:32 13:33 13:34 13:35 13:35 13:35 13:35 13:35 13:35 13:35 13:35 13:35 13:35 13:35 13:35 13:35 13:35

Ratio R=0.2/R=0.3 consistent with vacuum jets

LHC: Di-jet asymmetry E _{*T* 1}− E _{*T* 2} Dijet imbalance quantified by asymmetry variable AJ $A_J \equiv$ **Towers** E_{T1} + E_{T2} $\begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 &$ 35-280 GeV) embedded to HIJING with no quenching with no question of the second with no question of the second with no questions. Warning: Aj is sensitive to background fluctuations! Sub-leading jet in opposite hemisphere. Highest ET jet in opposite Need proper treatment in the data. Events selected 10-20% 20-40%

remember the 2-hadron correlations...

3.2 Track-jet correlations **17** CMS - quantifying the di-jet

asymmetry

Phe fractional imbalance: with vertical bars are shown as circles with vertical bars and brack-

- grows with collision centrality and reaches a much larger value than in PYTHIA or PYTHIA+DATA
- sharehous integrates for the birde states integrals cancel in the state set.
- clearly visible even for the highest-pT jets observed in the data set
- shown yishold even for the inghest-pp jets observed in the data set.
The nt + dependence of the evcess imbalance is compatible with either a constant difference or a constant fraction. eye is drawn at the value for pure PYTHIA for the lowest *p*^T bin. - the $p_{T,1}$ dependence of the excess imbalance is compatible with either a constant difference or a constant fraction of pT_1 .

di-jet asymmetry: where does the

energy go?

The momentum difference

 ad by low pT particle balanced by low-pT particles $\overline{}$ $\frac{1}{2}$

di-jet asymmetry: where does the

energy go?

 $p_{T,2}$ > 30 GeV/c

man function of trigger jet (1) pr $so GeV/c$ concertive of the city of 10^{10}

 \overline{C} collision central to collision collisions, corresponding to selections, corresponding to selections of \overline{C} Ratio follows the PYTHIA+HYDJET reference with the as points with vertical bars and brackets indicating the statistical and systematic uncertainties, \mathbf{r} same rate - constant offset over 200 GeV in pT

Modified jet fragmentation - an expectation from jet quenching $\xi = ln(E_{jet}/p_{hadron})$ $\frac{\text{d}N^\text{h}}{\text{d}\xi}(\xi,\tau)$ LEP Data Shown Jet vector • OPAL, $\sqrt{s} = 192 - 209$ GeV 14 $\overline{\text{in vacuum}}$, $E_{jet} = 100 \text{ GeV}$ ---- in medium, E_{jet} =100 GeV Jet quenching 10 8 "hump-back plateau" 6 projection 4 p_T ^{hadron_{~2} GeV} hadron $\overline{2}$ $\frac{1}{6}$ ξ =ln(E_{Jet}/p_{hadron}) 3 $\overline{2}$ High momentum Low momentum hadrons hadrons

in Papier Road

Jet fragmentation in Heavymental uncertainties, to that of jets produced in *pp* collisions as shown in Fig. 9 [132]. This suggests that \mathcal{L} and the additional energy radiated by the matter, and the fraction of the fragments of the model o calculations and intelligence of the combine elastic and inelastic and inelastic parton elastic parton elastic

The distribution of particle momenta inside jets normalized to the jets normalized to the same, within experimental $\mathcal{L}_\mathbf{z}$

ion collisions of radiated gluons by the medium, have been able to reproduce the increased energy asymmetry of different \mathcal{L}

CMS observation: Fragmentation of jets that lost energy consistent with jet fragmentation in proton-proton (vacuum) - similar observations by ATLAS

Transverse jet structure:

sition of *Talks: M. Veldhoen, J F Grosse-Oetringhaus* Internal composition of HI jets: 119

proton/pion ratio within a jet

- · No evidence of medium-induced modification
- Caution: physics evolves rapidly with pT in this region RAA at high-pT si. p_{Ttrig} :=> fragmentation bias
to all

 $(RHIC8LHC)$ Note: consistent with RAA at high-pT - similar to all species

Photon-jet

121 Direct photon (jet) measurement

An experimental chart... of the effort(!)

Direct photon(-jet) measurement ¹²²

pp collisions at 7 TeV [33] at about the 10% level in the region D*f >* 2*p*/3. The result that s esure: Photomation fetation observation of an unmodified D*f* correlation in dijet events [10]. **3 Results** 3.31 \mathbf{z} **3** \mathbf{z} \mathbf{z} \mathbf{z} \mathbf{z} \mathbf{z} \mathbf{z} \mathbf{z} \mathbf{z} Result: Photon(ΔE=0)-jet(^ΔE>0)

The asymmetry ratio $x_{J\gamma} = p_T^{\text{jet}}/p_T^{\gamma}$ is used to quantify the photon+jet momentum imbalance.

Summa summarum

- High energy heavy-ion collisions: Hot and dense (opaque to high-energy partons) quarkgluon plasma
- Hadron spectra suppressed (both at RHIC and LHC); Correlations of hadrons (proxies for 2->2 jet process) consistent with jet quenching
- Fully reconstructed jets suppressed (pT dependence of the suppression pattern different than for hadrons) -> constant fractional energy loss (?) -> up to highest jet energies measured (RHIC & LHC)
- The observed jets consistent with unmodified (vacuum?) fragmentation (within the current experimental assessment); The radiated energy "recovered" at large angles wrt jet axis
	- Also: no indication for particle type composition (p/pion etc) modifications of high-pT jets
	- Similar to jet-jet, the photon-jet correlations do NOT show de-correlation beyond p-p case (recoil jet also with unmodified fragmentation) Check the extra slides for more...
- Models explaining the phenomena being put forward.

RHIC jet results and examples of other observables (correlations) from LHC...

Do we understand everything about jet quenching and what fully reconstructed jet observables tell us? NO! But we learned already a lot... and this is just a good beginning!

Always good to ask: What is next?

- Improved control of the jet reconstruction in HIC still improvements possible (less biases, other observables) conceptually different approaches in making...
- New observables? Hadron-jet etc; Rates for 2+1+1 events? Structure of the jet with improved low-pT resolution? (subjets ?)
- Correlation of jets with the "soft" background and other observables? (low/intermediate-pT hadron correlations - take a look at the extra slides...) Look at the extra slides!
- Heavy-flavor jets? and their correlations?
- Energy-evolution of jet quenching more to learn? Higher energy (LHC)... RHIC still working on jets! Various collision systems...

...worth to look forward to... Your ideas can make a difference!

References (and refs therein!)

- Jet reconstruction (p-p and HIC), algorithms etc FastJet : http://fastjet.fr/about.html
- PHENIX results: http://www.phenix.bnl.gov/results.html
- STAR results: http://drupal.star.bnl.gov/STAR/publications
- ALICE results:<http://aliceinfo.cern.ch/ArtSubmission/publications>
- ATLAS HI results: https://twiki.cern.ch/twiki/bin/view/AtlasPublic/HeavyIonsPublicResults
- CMS HI results:<http://cms.web.cern.ch/org/cms-papers-and-results>
- Overview of first LHC results: Mueller, Wysloluch, Schuckraft: http://arxiv.org/abs/1202.3233
- Hard Probes 2012 conference:
	- <http://agenda.infn.it/conferenceOtherViews.py?view=standard&confId=4157>

Extra slides

• Did not fit for time reasons but also relevant(!)... make sure you go through these as well.

QCD Thermodynamics -

calculation

Particle detection

Glauber Monte Carlo

- Glauber model: geometrical picture of AA collision
	- Straight-line nucleon trajectories
	- N-N cross-section independent of the number of collisions the nucleons have undergone before

Nuclear density profile: Woods-Saxon (2pF)

$$
\rho(r) = \rho_o \cdot \frac{1}{1 + \exp\left(\frac{r - R}{d}\right)}
$$

- Radius=6.62±0.06fm
- $-$ skin depth=0.546 \pm 0.01fm
- $-$ Intra-nucleon distance=0.4 \pm 0.4fm
- **Nucleon-Nucleon inelastic cross section** σ_{NN} =64±5 mb at 2.76 TeV
- Estimate uncertainty by varying model **assumptions**

Particle multiplicity & centrality

dN/dη scales faster than pp

- Trend predicted by some saturation model
- Excellent agreement with LHC experiments
- Energy density $\times \tau_0 \approx 3 \times R$ HIC

$$
\varepsilon \geq \frac{dE_T/d\eta}{\tau_0 \pi R^2} = \frac{3}{2} \langle E_T/N \rangle \frac{dN_{\rm ch}/d\eta}{\tau_0 \pi R^2}
$$

Scaling similar to RHIC:

• Contribution of hard processes $(N_{coll}$ scaling)?

Classes of models

- Saturation
- 2 components (hard/soft)
- \rightarrow models incorporating moderation of multiplicity (shadowing/saturation) favoured

• More on two-particle correlations

Sensitivity of particle correlations to different underlying physics

"Beyond" v2

higher moments -> fluctuations / hotspots

Single event!

Higher harmonics - measured

We observe significant v₃ which compared to v_2 has a different centrality dependence $|$ v₃ - triangular flow : $|$ - weak centrality dependence - vanishes as expected when \parallel measured w.r.t. reaction plane The centrality dependence and
magnitude are similar to

● Weak centrality dependence Similar pT dependence for all v_n

rhor hormonics additional $\cos \theta$ constraints on n/s Higher harmonics - additional constraints on η/s

● η/s small, similar as at RHIC

● Higher harmonics provide $\frac{1}{2}$ constraints on the constraints on α η/s small, similar as at RHIC

Two-particle correlations -

Fourier decomposition Fourier decomposition 22 Fourier decomposition 22 trier decomposition

Correlations & hydrodynamics...

 $\| \cdot \|$ ong range correlations – collective flow; the coefficients must factorize such that; $\| \cdot \|$ Long range correlations – collective flow: the coefficients must factorize such that: rollective flow: the coefficients must factorize such that:

 $\mathbb{R}^{\mathbb{Z}}$

Global fits show: Show:

- **Collective flow dominates to about 3-4 GeV/c for all n>1** \mathbf{c}
- **Description breaks for high pT or peripheral collisions**
- **For low pT: double peak and ridge structures seen in two particle correlations are** naturally explained by measured anisotropic flow coefficients

are naturally explained by measured anisotropic flow coefficients α

Jet-medium-flow coupling τ armore that τ are the salgador of τ *Measuring the Collective Flow with Jets*

via two particle correlations?

STAR preliminary • N. Armesto, C. Salgado, U. Wiedemann: *Measuring the Collective Flow with Jets*

Δη **rms** [PRL 93,242301 (2004)]

Jet-peak shape

evolution - intermediate pT

Unique Peripheral a Peripheral and p-p similar shape • **Peripheral and pp similar** • **Strong** *p***T dependence** Strong pT dependence Wider peak in central collisions => Characterize the peak

b) η -gap subtracted

 $\overline{}$

m condthes of the correlations in Measuring widths of the correlations in

Talks: A. Morsch, J F Grosse-Oetringhaus azimuth and pseudo-rapidity

 $2 < p_{T,t} < 3$ 1 $< p_{T,a} < 2$ GeV/c \sum_{p}^{100} (fit)
 \sum_{p}^{100} (fit)
 \sum_{p}^{100} (s) $4 < p_{T,t} < 8$ 2 $< p_{T,a} < 3$ GeV/c 1.8 Vacuum Static medium: Flowing medium: ៖
[៖]1.6 PRELIMINARY Anisotropic shape (reference) Broadening $\sigma_{\Delta n}$ $1.4¹$ $\sigma_{\Delta\phi}^{}$ 1.2 $\sigma_{\!A\phi}$ 0.5 1.0 0.8 0.4 0.6 0.4 0.2 0.2 20 40 60 80 pp 0.0 Centrality $_{\text{o.o.}}$ Centrality, $_{\text{o}}$ 100 = pp $\Delta \omega$ (rad.). An **data points:** $\mathcal{L} = \mathcal{L} = \mathcal$

Hard Probes 2012, Canada Probes

Mea: $\frac{1}{2}$ interesting $\frac{1}{2}$ radiation $\|$ longitudinal flow (?) rveas*u*i Measure of jets interactions with

pT in this region is a set of the contract of the contract of

- AMPT (**A M**ulti**P**hase **T**ransport Code)
	- Initial conditions simulated using HIJING
	- Parton scattering
	- Hadronization: Lund model + coalescence
	- Hadron scattering
- AMPT describes the main features of the near-side shape evolution observed in data

• Jets at RHIC...

Jets at RHIC in HIC Suppression without de-correlation in *Cu*+*Cu*

Jets at RHIC in HIC Tots of RUTC in UTC

Work on final results in progress... Work on final results in progress...

Recoil jet spectrum at RHIC

Trigger-jet: biased towards surface

- strong fragmentation bias ~ vacuum jet

• Selecting biased trigger jet maximizes path length for the recoil (b-2-b) jets: extreme selection of jet population • Significant suppression in di-jet coincidence measurements!

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RHIC: Jet-hadron coincidences

suppression is compensated for by low- p_T^{assoc} enhancement.

Reminder on fragmentation bias...

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• Fragmentation bias! - nature is kind and (in most cases) will give you what you ask for - perhaps NOT what you WANT

Thank you!

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- For graphics/slides from: B. Cole, P. Govoni, M. Nguyen, T. Hemmick, P. Jacobs, M. Floris, M. van Leeuwen, C. Loizides, A. Morsch, J. Putschke, C. Roland, M. Rybář, G. Salam, Y. Shi Lai, G. Soyez, I. Wingerter
- For the material by collaborations: ALICE, ATLAS, CMS, PHENIX, STAR

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