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The lightcone
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Kernel and
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Large c limits

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McGill University

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Quantum Regge Trajectories and the Virasoro Analytic Bootstrap

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Conformal block decomposition

All CFTs have OPE (here scalar)

$$\phi(x)\phi(0) = \sum_{\mathcal{O}} f_{\phi\phi\mathcal{O}} C(x, \partial)\mathcal{O}(0)$$

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Conformal block decomposition

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$$\phi(x)\phi(0) = \sum_{\mathcal{O}} f_{\phi\phi\mathcal{O}} C(x, \partial)\mathcal{O}(0)$$

Consider using it for 12 and 34 (s-channel) in $d \geq 3$

$$\langle \phi(x_1)\phi(x_2)\phi(x_3)\phi(x_4) \rangle = \frac{\sum_{\mathcal{O}} f_{\phi\phi\mathcal{O}}^2 G_{\Delta_{\mathcal{O}}, \ell_{\mathcal{O}}}^{\Delta_{\phi}}(z, \bar{z})}{(x_{12})^{2\Delta_{\phi}}(x_{34})^{2\Delta_{\phi}}}$$

with $G_{\Delta_{\mathcal{O}}, \ell_{\mathcal{O}}}^{\Delta_{\phi}}$ conformal blocks and z, \bar{z} conformal cross-ratios

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with $G_{\Delta_{\mathcal{O}}, \ell_{\mathcal{O}}}^{\Delta_{\phi}}$ conformal blocks and z, \bar{z} conformal cross-ratios

Write sum in terms of twist $\tau = \Delta - \ell$

Crossing symmetry

Can do 14 and 23 instead (t-channel) and get same thing

$$\sum_{\mathcal{O}} f_{\phi\phi\mathcal{O}}^2 G_{\tau,\ell}^{\Delta_\phi}(z, \bar{z}) = \left(\frac{z\bar{z}}{(1-z)(1-\bar{z})} \right)^{\Delta_\phi} \sum_{\mathcal{O}'} f_{\phi\phi\mathcal{O}'}^2 G_{\tau',\ell'}^{\Delta_\phi}(1-z, 1-\bar{z})$$

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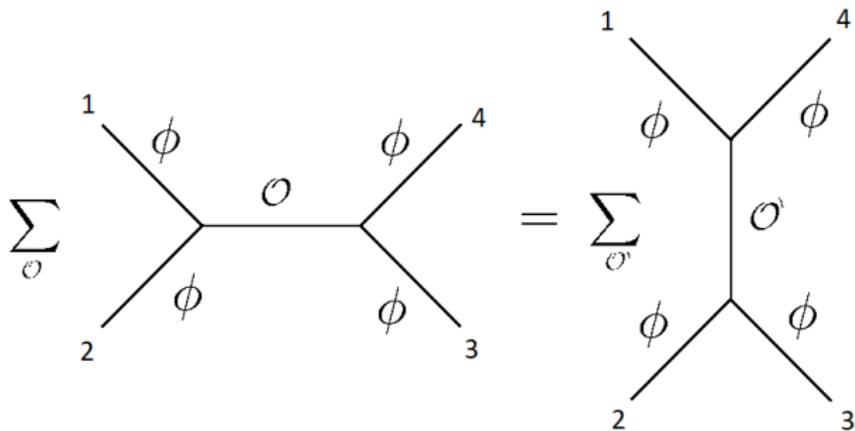
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Crossing symmetry

Can do 14 and 23 instead (t-channel) and get same thing

$$\sum_{\mathcal{O}} f_{\phi\phi\mathcal{O}}^2 G_{\tau,\ell}^{\Delta_\phi}(z, \bar{z}) =$$

$$\left(\frac{z\bar{z}}{(1-z)(1-\bar{z})} \right)^{\Delta_\phi} \sum_{\mathcal{O}' } f_{\phi\phi\mathcal{O}'}^2 G_{\tau',\ell'}^{\Delta_\phi}(1-z, 1-\bar{z})$$



Lightcone limit

Take $\bar{z} \rightarrow 1$, t-channel blocks behave as

$$G_{\tau', \ell'}^{\Delta, \phi}(1-z, 1-\bar{z}) \approx (1-\bar{z})^{\frac{\tau'}{2}} K_{\Delta'+\ell'}(1-z)$$

\Rightarrow t-channel dominated by identity!

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Lightcone limit

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Further take $z \rightarrow 0$, s-channel blocks behave as

$$G_{\tau, \ell}^{\Delta_\phi}(z, \bar{z}) \approx z^{\frac{\tau}{2}} \log(1-\bar{z})$$

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$$G_{\tau, \ell}^{\Delta_\phi}(z, \bar{z}) \approx z^{\frac{\tau}{2}} \log(1-\bar{z})$$

Crossing symmetry becomes

$$\sum_{\tau, \ell} f_{\phi\phi\mathcal{O}}^2 z^{\frac{\tau}{2}} \log(1-\bar{z}) = \frac{z^{\Delta_\phi}}{(1-\bar{z})^{\Delta_\phi}} + \dots$$

Double twist operators

Impossible to reproduce t-channel singularity with finite number of terms

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Impossible to reproduce t-channel singularity with finite number of terms

\Rightarrow Need infinite family of operators with

$$\tau = 2\Delta_\phi + 2n$$

for $\ell \rightarrow \infty$

Double twist operators

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Impossible to reproduce t-channel singularity with finite number of terms

⇒ Need infinite family of operators with

$$\tau = 2\Delta_\phi + 2n$$

for $\ell \rightarrow \infty$

Call these operators “double twist”, schematically

$$[\phi\phi]_{n,\ell} = \phi \square^n \partial^\ell \phi$$

Double twist operators

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Explicitly inverting crossing gives the OPE coefficients

Mean Field Theory

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t-channel identity \Rightarrow s-channel "double twists"

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t-channel identity \Rightarrow s-channel "double twists"

Reproduces Mean Field Theory: CFT with correlators given by Wick contractions, contain only double twist operators

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RESULT: Every CFT behaves as MFT at large spin

Mean Field Theory

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t-channel identity \Rightarrow s-channel "double twists"

Reproduces Mean Field Theory: CFT with correlators given by Wick contractions, contain only double twist operators

RESULT: Every CFT behaves as MFT at large spin

Including subleading operators in t-channel gives corrections to OPE and anomalous dimensions

$$\gamma_{n,\ell} \sim \frac{1}{\ell^\tau}$$

Regge trajectories

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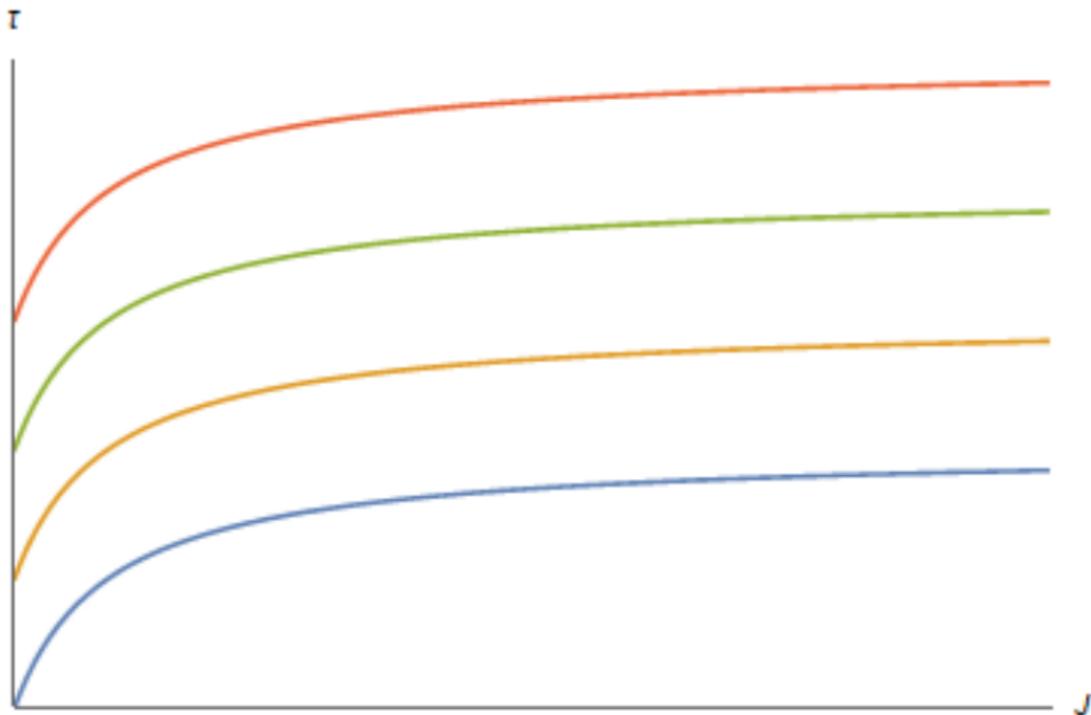
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Inversion formula

Can write 4-point function as

$$\langle \phi(x_1)\phi(x_2)\phi(x_3)\phi(x_4) \rangle \sim \sum_{\ell=0}^{\infty} \int_{\frac{d}{2}-i\infty}^{\frac{d}{2}+i\infty} d\Delta C(\Delta, \ell) G_{\Delta, \ell}(z, \bar{z})$$

where C has poles at physical operator with residues giving the OPE coefficients

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Inversion formula

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Can write 4-point function as

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where C has poles at physical operator with residues giving the OPE coefficients

Simon's formula inverts this

$$C(\Delta, \ell) \propto \int_0^1 \int_0^1 dz d\bar{z} M_{\Delta, \ell}(z, \bar{z}) \text{dDisc}[\langle \phi(x_1)\phi(x_2)\phi(x_3)\phi(x_4) \rangle]$$

$6j$ symbols

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Inserting identity in inversion formula gives MFT result

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Inversion of single block = $6j$ symbol

$6j$ symbols

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Inserting identity in inversion formula gives MFT result

Inserting other operators gives corrections

Inversion of single block = $6j$ symbol

\Rightarrow $6j$ symbols rewrite t-channel data into s-channel data

Problems in 2d

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What is wrong in $2d$?

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What is wrong in 2d?

- Virasoro blocks not known

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What is wrong in $2d$?

- Virasoro blocks not known
- No twist gap (T , T^2 , etc. have zero twist)

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FKW already studied this in large c limit for HHLL with Virasoro vacuum block

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We will take finite c and reproduce their results.

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2d CFT

Conformal transformations factorize into holomorphic and anti-holomorphic

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Conformal transformations factorize into holomorphic and anti-holomorphic

\Rightarrow Conformal blocks factorize

$$G(z, \bar{z}) = \mathcal{F}(h|z)\bar{\mathcal{F}}(\bar{h}|\bar{z})$$

with $h = \Delta + \ell$ and $\bar{h} = \Delta - \ell$

2d CFT

Conformal transformations factorize into holomorphic and anti-holomorphic

⇒ Conformal blocks factorize

$$G(z, \bar{z}) = \mathcal{F}(h|z)\bar{\mathcal{F}}(\bar{h}|\bar{z})$$

with $h = \Delta + \ell$ and $\bar{h} = \Delta - \ell$

Crossing symmetry for $\langle \mathcal{O}_1(0)\mathcal{O}_2(z, \bar{z})\mathcal{O}_2(1)\mathcal{O}_1(\infty) \rangle$ is now

$$\sum_s (f_{12s})^2 \mathcal{F}_S(h_s, z) \bar{\mathcal{F}}_S(\bar{h}_s, \bar{z}) = \sum_t f_{11t} f_{22t} \mathcal{F}_T(h_t, 1-z) \bar{\mathcal{F}}_T(\bar{h}_t, 1-\bar{z})$$

Liouville notation

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Need to use new notation:

$$c = 1 + 6Q^2 \quad , \quad Q = b + b^{-1} \quad , \quad h = \alpha(Q - \alpha)$$

$$(h, c) \Rightarrow (\alpha, b)$$

Liouville notation

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Need to use new notation:

$$c = 1 + 6Q^2 \quad , \quad Q = b + b^{-1} \quad , \quad h = \alpha(Q - \alpha)$$

$$(h, c) \Rightarrow (\alpha, b)$$

Operators separate in two ranges:

$$\text{Discrete: } 0 < h < \frac{c-1}{24} \longleftrightarrow 0 < \alpha < \frac{Q}{2}$$

$$\text{Continuum: } h \geq \frac{c-1}{24} \longleftrightarrow \alpha = \frac{Q}{2} + iP$$

Definition of the kernel

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Rewrite t-channel (holomorphic) Virasoro blocks into s-channel blocks

$$\mathcal{F}_T(\alpha_t, 1-z) = \int_C \frac{d\alpha_s}{2i} \mathbb{S}_{\alpha_s \alpha_t} \mathcal{F}_S(\alpha_s, z)$$

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Impressive that it is known since blocks themselves not known

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Impressive that it is known since blocks themselves not known

Poles at $\alpha_s = \alpha_1 + \alpha_2 + mb + nb^{-1}$ and reflexions $\alpha \rightarrow Q - \alpha$

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Impressive that it is known since blocks themselves not known

Poles at $\alpha_s = \alpha_1 + \alpha_2 + mb + nb^{-1}$ and reflexions $\alpha \rightarrow Q - \alpha$

- For $\alpha_t = 0$, single poles
- For $\alpha_t \neq 0$, double poles

Analytic structure

When $\alpha_1 + \alpha_2 > \frac{Q}{2}$, C is simple

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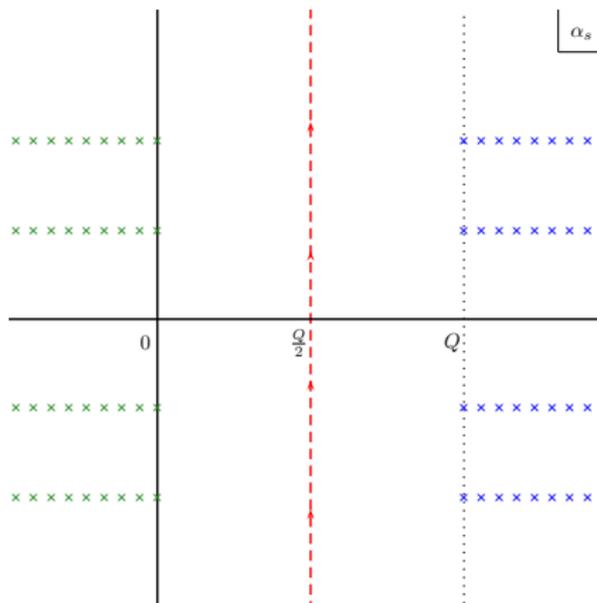
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Analytic structure

When $\alpha_1 + \alpha_2 > \frac{Q}{2}$, C is simple



Analytic structure

When $\alpha_1 + \alpha_2 < \frac{Q}{2}$, poles at $\alpha_m = \alpha_1 + \alpha_2 + mb$ can cross axis

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Analytic structure

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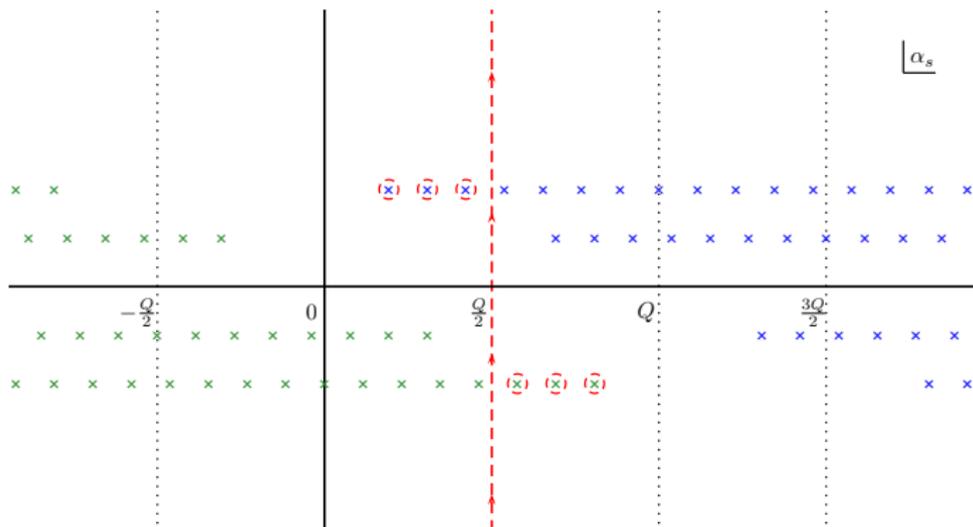
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When $\alpha_1 + \alpha_2 < \frac{Q}{2}$, poles at $\alpha_m = \alpha_1 + \alpha_2 + mb$ can cross axis



Support of the kernel

- For $\alpha_1 + \alpha_2 > \frac{Q}{2}$,

$$\mathcal{F}_T(\alpha_t) = \int_0^\infty dP \mathbb{S}_{\alpha_s \alpha_t} \mathcal{F}_S \left(\alpha_s = \frac{Q}{2} + iP \right)$$

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Support of the kernel

- For $\alpha_1 + \alpha_2 > \frac{Q}{2}$,

$$\mathcal{F}_T(\alpha_t) = \int_0^\infty dP \mathbb{S}_{\alpha_s \alpha_t} \mathcal{F}_S \left(\alpha_s = \frac{Q}{2} + iP \right)$$

- For $\alpha_1 + \alpha_2 < \frac{Q}{2}$,

$$\begin{aligned} \mathcal{F}_T(\alpha_t) = & -2\pi \sum_m \operatorname{Res}_{\alpha_s = \alpha_m} \{ \mathbb{S}_{\alpha_s \alpha_t} \mathcal{F}_S(\alpha_s) \} \\ & + \int_0^\infty dP \mathbb{S}_{\alpha_s \alpha_t} \mathcal{F}_S \left(\alpha_s = \frac{Q}{2} + iP \right) \end{aligned}$$

with sum over $\alpha_m < \frac{Q}{2}$

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Crossing with fusion

Rewrite t-channel into s-channel with kernel tells us what must be there in the s-channel to reproduce what appears in t-channel.

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Crossing with fusion

Rewrite t-channel into s-channel with kernel tells us what must be there in the s-channel to reproduce what appears in t-channel.

Consider $\alpha_1 + \alpha_2 < \frac{Q}{2}$ and $\bar{\alpha}_1 + \bar{\alpha}_2 > \frac{Q}{2}$ and
individual t-channel exchange

$$\int_{\mathcal{I}} d\alpha_s d\bar{\alpha}_s \rho_{12s} \mathcal{F}_S(\alpha_s) \bar{\mathcal{F}}_S(\bar{\alpha}_s) = \int_0^\infty d\bar{P} \bar{\mathbb{S}}_{\bar{\alpha}_s \bar{\alpha}_t} \bar{\mathcal{F}}_S \left(\bar{\alpha}_s = \frac{Q}{2} + i\bar{P} \right)$$
$$f_{11t} f_{22t} \left[-2\pi \sum_m \text{Res}_{\alpha_s = \alpha_m} \{ \mathbb{S}_{\alpha_s \alpha_t} \mathcal{F}_S(\alpha_s) \} + \int_0^\infty dP \mathbb{S}_{\alpha_s \alpha_t} \mathcal{F}_S \left(\alpha_s = \frac{Q}{2} + iP \right) \right]$$

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What is needed to reproduce identity $\alpha_t = \bar{\alpha}_t = 0$?

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What is needed to reproduce identity $\alpha_t = \bar{\alpha}_t = 0$?

- 1 Family of operators with $\alpha = \alpha_m < \frac{Q}{2}$ (in discrete spectrum) for each $\bar{\alpha}$ in continuum \Rightarrow “Quantum” Regge trajectories
- 2 Operators with α and $\bar{\alpha}$ in continuum

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OPE coefficients of Regge operators given by

$$\rho_{12m} = -2\pi \bar{S}_{\bar{\alpha}_s \mathbb{I}} \operatorname{Res}_{\alpha_s = \alpha_m} S_{\alpha_s \mathbb{I}}$$

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OPE coefficients of Regge operators given by

$$\rho_{12m} = -2\pi \bar{S}_{\bar{\alpha}_s \mathbb{I}} \operatorname{Res}_{\alpha_s = \alpha_m} S_{\alpha_s \mathbb{I}}$$

This is called Virasoro Mean Field Theory!

Corrections

Assume other operators give small corrections

$$\begin{aligned}(\rho_{12m} + \delta\rho_{12m})\mathcal{F}_S(\alpha_m + \delta\alpha_m)\bar{\mathcal{F}}_S \approx & \bar{\mathcal{F}}_S (\rho_{12m}\mathcal{F}_S(\alpha_m) + \\ & \delta\rho_{12m}\mathcal{F}_S(\alpha_m) + \\ & \rho_{12m}\delta\alpha_m\partial\mathcal{F}_S(\alpha_m))\end{aligned}$$

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Corrections

Assume other operators give small corrections

$$(\rho_{12m} + \delta\rho_{12m})\mathcal{F}_S(\alpha_m + \delta\alpha_m)\bar{\mathcal{F}}_S \approx \bar{\mathcal{F}}_S (\rho_{12m}\mathcal{F}_S(\alpha_m) + \delta\rho_{12m}\mathcal{F}_S(\alpha_m) + \rho_{12m}\delta\alpha_m\partial\mathcal{F}_S(\alpha_m))$$

This leads to

$$\delta\alpha_m = f_{11t}f_{22t} \frac{\bar{\mathbb{S}}_{\bar{\alpha}_s\bar{\alpha}_t} \text{dRes}_{\alpha_s=\alpha_m} \mathbb{S}_{\alpha_s\alpha_t}}{\bar{\mathbb{S}}_{\bar{\alpha}_s\mathbb{I}} \text{Res}_{\alpha_s=\alpha_m} \mathbb{S}_{\alpha_s\mathbb{I}}}$$

$$\delta\rho_{12m} = -2\pi f_{11t}f_{22t} \bar{\mathbb{S}}_{\bar{\alpha}_s\bar{\alpha}_t} \text{Res}_{\alpha_s=\alpha_m} \mathbb{S}_{\alpha_s\alpha_t}$$

where dRes means the coefficient of double pole

Why dRes?

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Large c limits

Taylor expanding double pole at $x = x_0$ gives

$$s(x)f(x) = \left(\frac{d\text{Res}(s)}{(x - x_0)^2} + \frac{\text{Res}(s)}{x - x_0} + s(x_0) \right) \times (f(x_0) + (x - x_0)f'(x_0))$$

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$$\begin{aligned} s(x)f(x) &= \left(\frac{d\text{Res}(s)}{(x-x_0)^2} + \frac{\text{Res}(s)}{x-x_0} + s(x_0) \right) \\ &\quad \times (f(x_0) + (x-x_0)f'(x_0)) \\ &= \frac{f(x_0) d\text{Res}(s)}{(x-x_0)^2} + \frac{f(x_0) \text{Res}(s) + f'(x_0) d\text{Res}(s)}{x-x_0} + f'(x_0) \text{Res}(s) + \dots \end{aligned}$$

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$$\Rightarrow \text{Res}(s f) = f(x_0) \text{Res}(s) + f'(x_0) d\text{Res}(s)$$

Large spin asymptotics

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At large $\bar{\alpha}_s$

$$\delta\alpha_m \sim e^{-2\pi\bar{\alpha}_t\sqrt{\ell_s}}$$

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At large $\bar{\alpha}_s$

$$\delta\alpha_m \sim e^{-2\pi\bar{\alpha}_t\sqrt{\ell_s}}$$

\Rightarrow identity dominates at large spin!

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At large $\bar{\alpha}_s$

$$\delta\alpha_m \sim e^{-2\pi\bar{\alpha}_t\sqrt{\bar{\ell}_s}}$$

\Rightarrow identity dominates at large spin!

Spectrum of Quantum Regge trajectories at large spin:

$$h_m = h_1 + h_2 + m - 2(\alpha_1 + mb)(\alpha_2 + mb) + m(m+1)b^2 + \delta h_m$$

Quantum Regge trajectories

The Virasoro fusion kernel and its applications

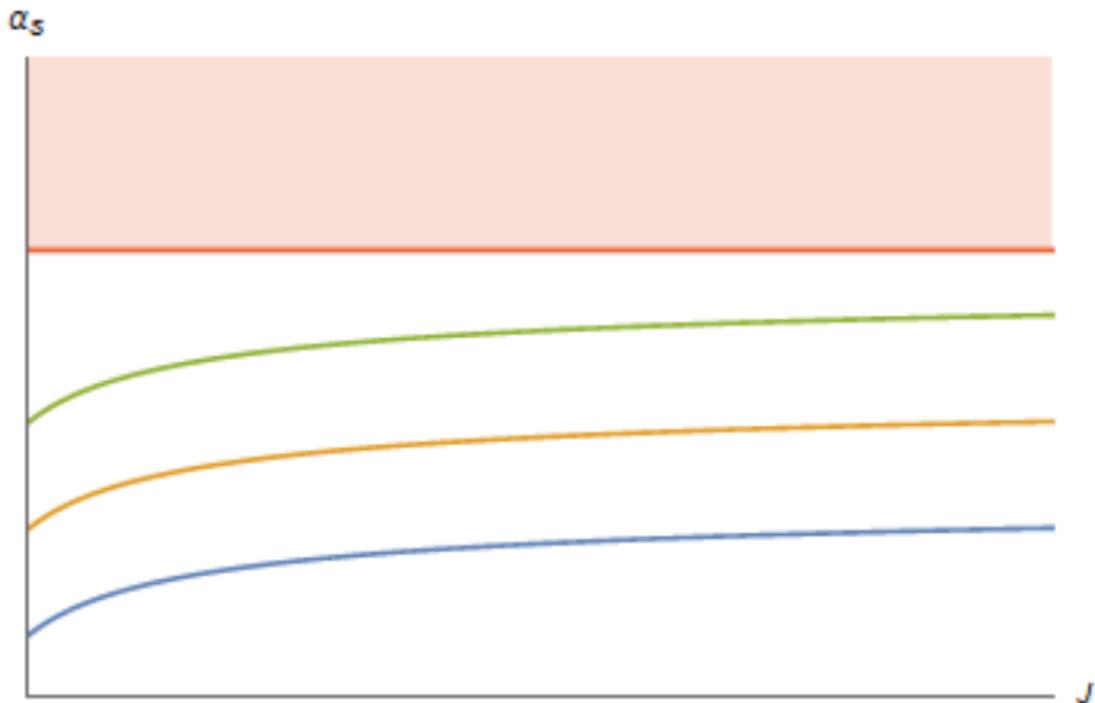
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Reproduce global results with $c \rightarrow \infty$ and h_i fixed

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Reproduce global results with $c \rightarrow \infty$ and h_i fixed

$$\Rightarrow \alpha = bh + O(b^3) \text{ as } b \rightarrow 0$$

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Reproduce global results with $c \rightarrow \infty$ and h_i fixed

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Infinite number of trajectories with

$$h_m = h_1 + h_2 + m + O(b^2)$$

Global limit

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Reproduce global results with $c \rightarrow \infty$ and h_i fixed

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Infinite number of trajectories with

$$h_m = h_1 + h_2 + m + O(b^2)$$

Checks:

- 1 Reproduce MFT from VMFT (exchange of identity)
- 2 Other t-channel reproduced
- 3 Next order in identity exchange gives T

Large c trajectories

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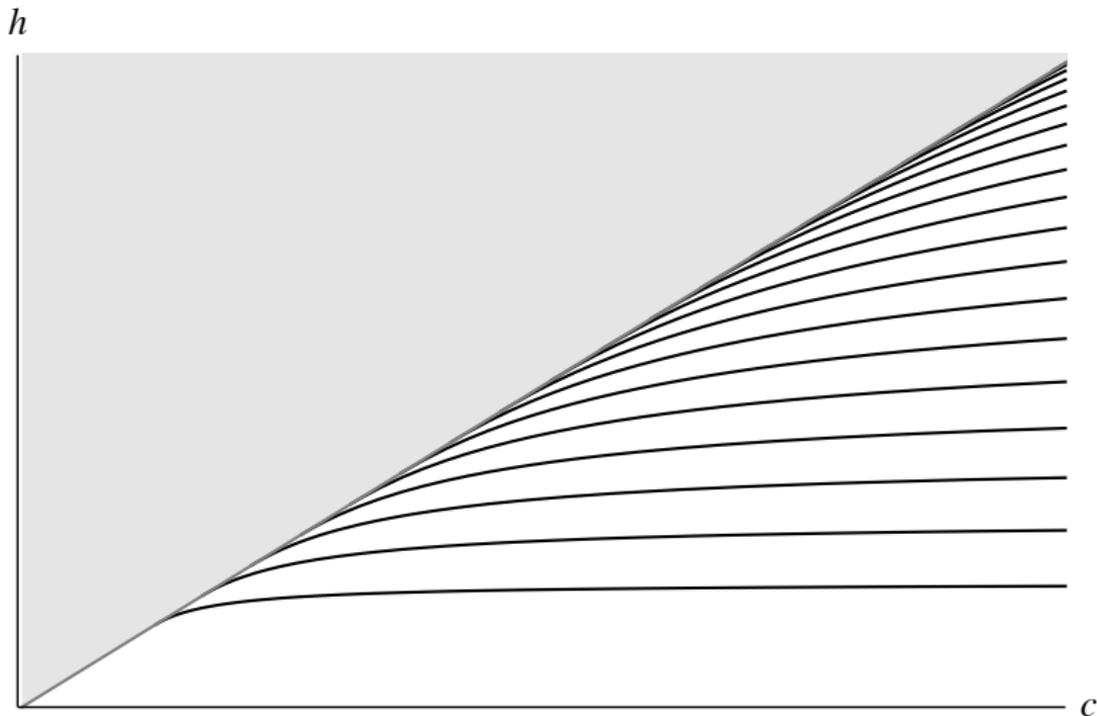
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Semiclassical limit

Again $c \rightarrow \infty$ but some operators heavy $h \sim c$

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Again $c \rightarrow \infty$ but some operators heavy $h \sim c$

$$\Rightarrow \alpha = \frac{Q}{2} + ib^{-1}p \text{ or } \alpha = \eta b^{-1} \text{ as } b \rightarrow 0$$

Semiclassical limit

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Large c limits

Again $c \rightarrow \infty$ but some operators heavy $h \sim c$

$$\Rightarrow \alpha = \frac{Q}{2} + ib^{-1}p \text{ or } \alpha = \eta b^{-1} \text{ as } b \rightarrow 0$$

When $m \ll b^{-1} \sim \sqrt{c}$, $h_1 = O(c) < \frac{c}{24}$ and $h_2 = O(1)$,
recover

$$h_m \approx h_1 + \sqrt{1 - \frac{24h_1}{c}}(h_2 + m)$$

same as FKW

Semiclassical limit

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same as FKW

When further take $\frac{h_1}{c} \ll 1$, recover

$$h_m \approx h_1 + h_2 - \frac{12h_1h_2}{c}$$

which can be derived from inversion formula

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Large c limits

- 1 Fusion kernel: write t-channel Virasoro block in terms of s-channel blocks

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- 1 Fusion kernel: write t-channel Virasoro block in terms of s-channel blocks
- 2 VMFT: inversion of identity Virasoro block

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- 4 Corrections to trajectories

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- 1 Fusion kernel: write t-channel Virasoro block in terms of s-channel blocks
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Other results

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Many other applications

- 1 Virasoro blocks at late time (information paradox)
- 2 Gravity interpretation
- 3 $z \rightarrow 1$ limit of Virasoro blocks
- 4 HHLL Virasoro blocks
- 5 2d lightcone bootstrap